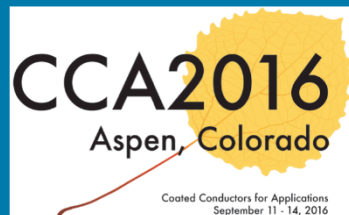


Angular dependence of the critical current density in coated conductors with an isotropic defect structure

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CCA, Aspen, Colorado, September 13th 2016

Outline

- Motivation
 - Angular dependence of critical currents in neutron irradiated coated conductors
- Neutron irradiation
 - Introduced defect structure
- Results
 - Anisotropic scaling approach (Ba-122 single crystals)
 - Resulting angular dependence of critical currents
 - Comparison with experimental results on coated conductors
- Conclusions



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Measurements of irradiated coated conductors were performed by **Michal Chudy**.





Famous view of Maroon Bells

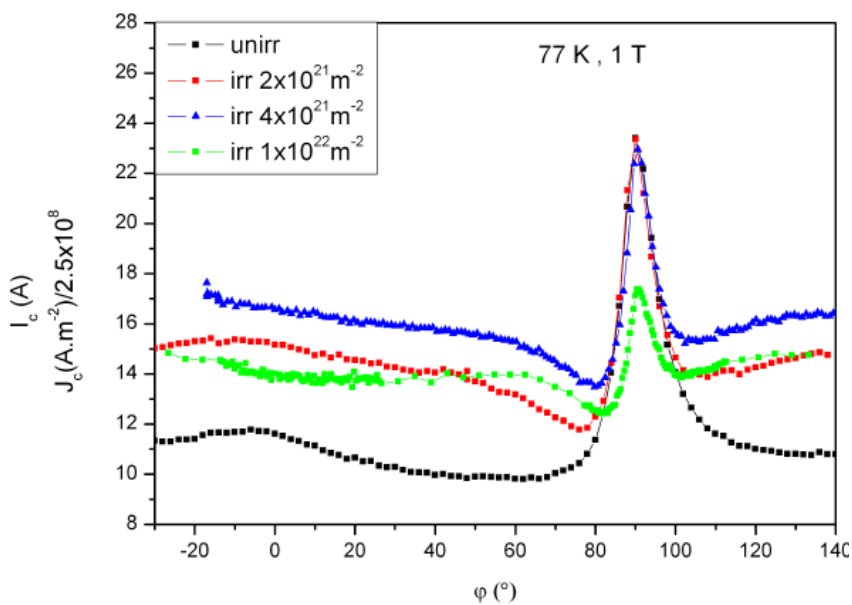


Change of perspective

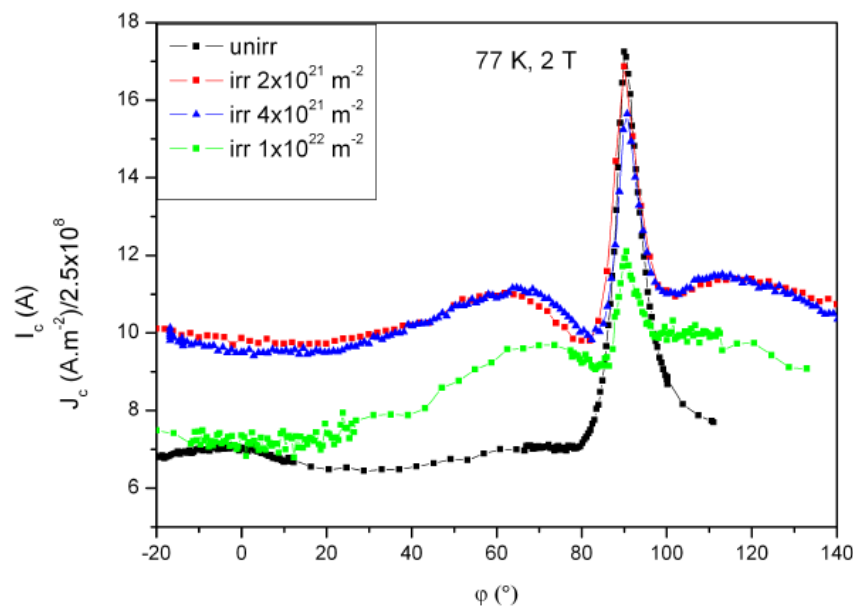


may be quite rewarding!

Neutron irradiated coated conductors



Peak at $H \parallel c$



Peak at intermediate angles

Resulting from isotropic defects?

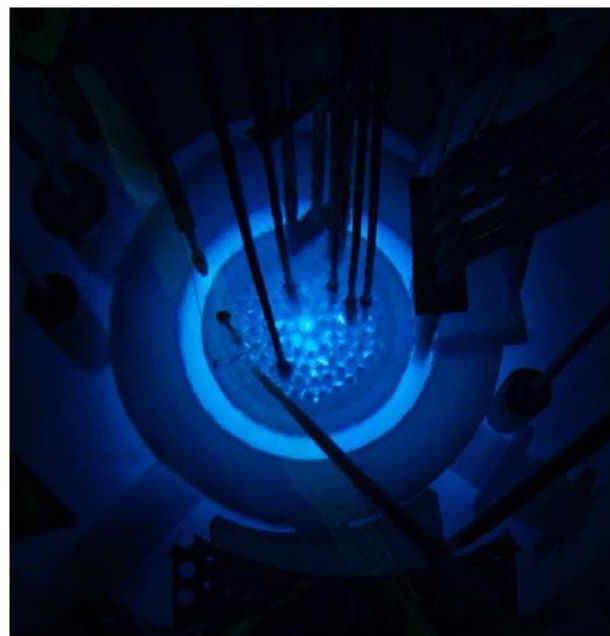


NEUTRON IRRADIATION



TRIGA MARK-II reactor

Fast (>0.1 MeV) neutron flux density: $4.5 \times 10^{16} \text{ m}^{-2}\text{s}^{-1}$



Neutron Irradiation: Created Defects

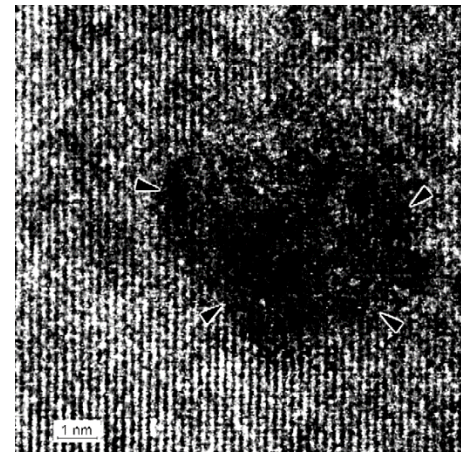
Direct collisions (high energy neutrons)

Defect cascades ($E > 0.1$ MeV)

$\varnothing \sim 5$ nm

Density: 10^{22} m⁻³ at a fluence of 2×10^{21} m⁻²

($d_{av} \sim 46$ nm, $B_f \sim 1$ T)



Smaller defects:

Single displaced atoms, clusters of point defects....

Randomly distributed, NOT correlated, nearly spherical



RESULTS



Anisotropic scaling approach

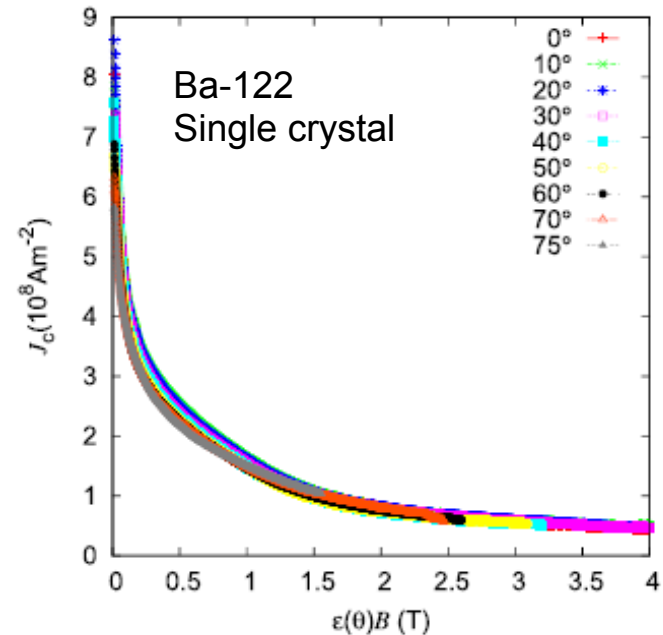
$$J_{\downarrow c}(B, \theta) = J_{\downarrow c}(B\epsilon(\theta))$$

$$\epsilon(\theta) = \sqrt{\gamma \uparrow - 2 \sin^2(\theta) + \cos^2(\theta)}$$

G. Blatter et al., Phys. Rev. Lett. **68** (1992) 875

Collective pinning theory, single vortex pinning regime (weak pinning, low fields)

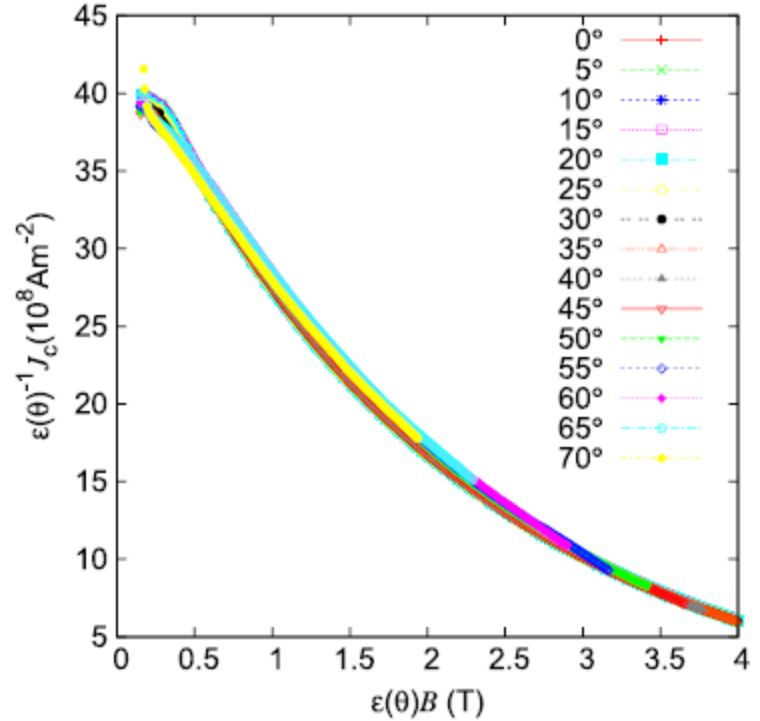
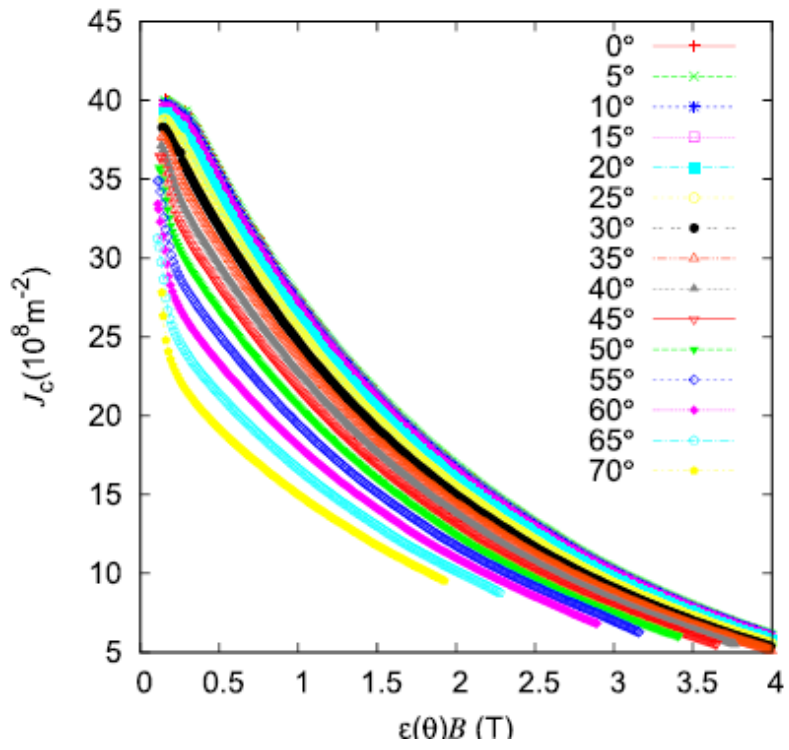
Angle-resolved magnetization measurements



V. Mishev et al.,
Supercond. Sci. Technol. **28** (2015) 102001C



Neutron irradiated Ba-122 (Co-doped)

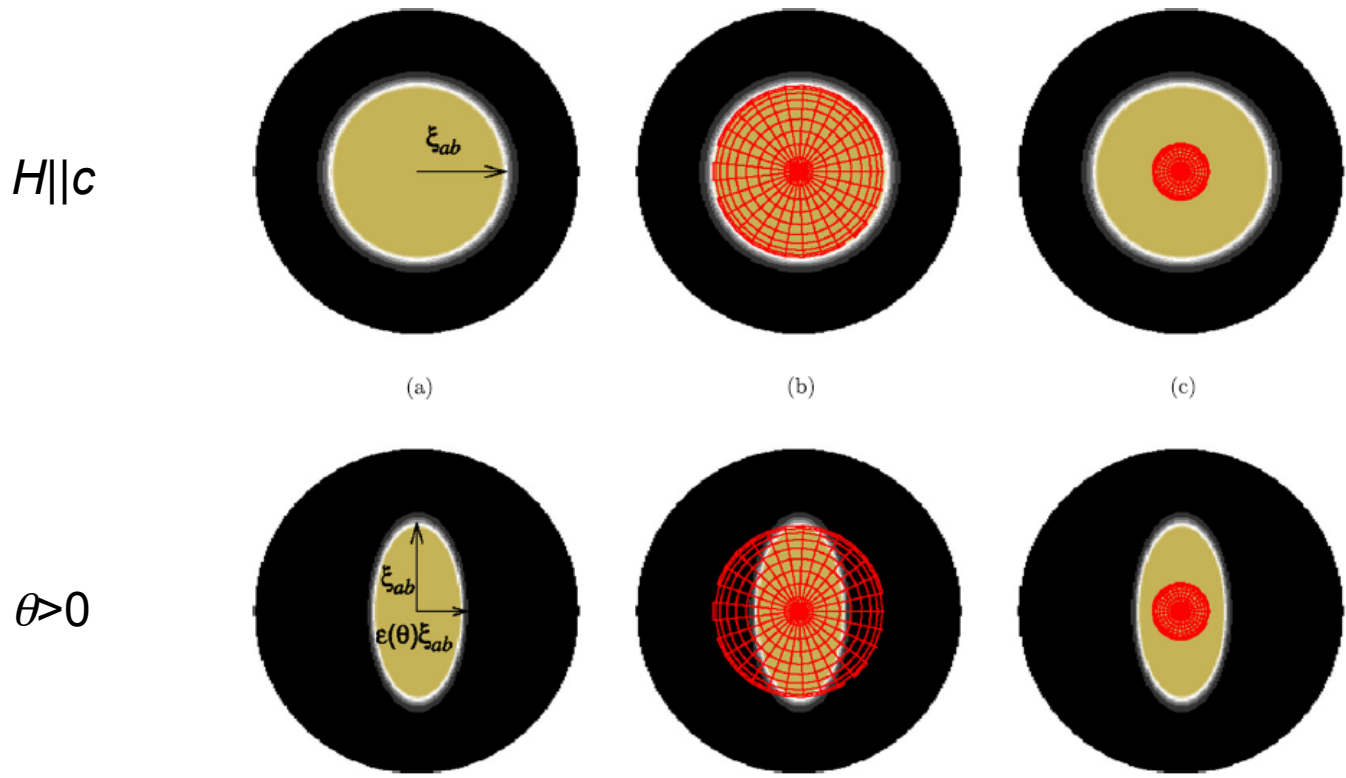


V. Mishev et al., Supercond. Sci. Technol. **28** (2015) 102001

Scaling of field **and** the critical current density is needed.



Point-like vs. spherical defects



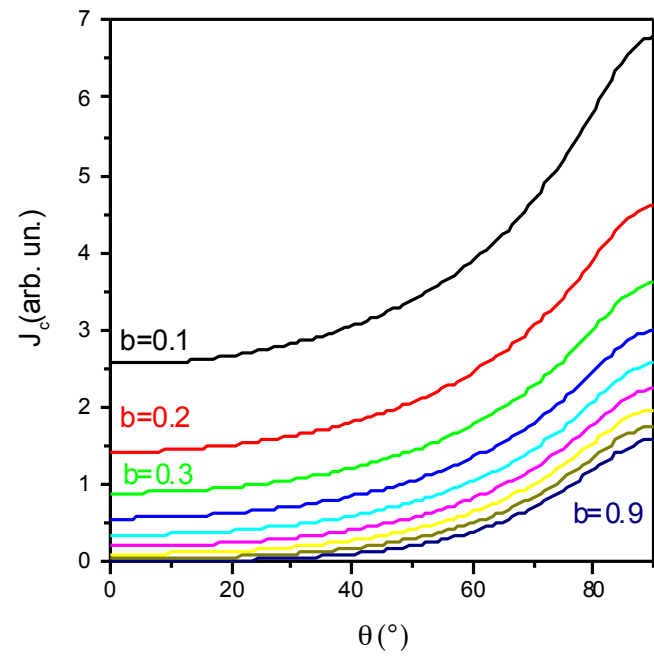
V. Mishev et al., Supercond. Sci. Technol. **28** (2015) 102001

(Anisotropic scaling approach: length has to be scaled)



Implications for angular dependence of J_c

$$J_c \propto B^{\gamma-0.5} (1-B/B_{irr})^{\gamma/2}$$

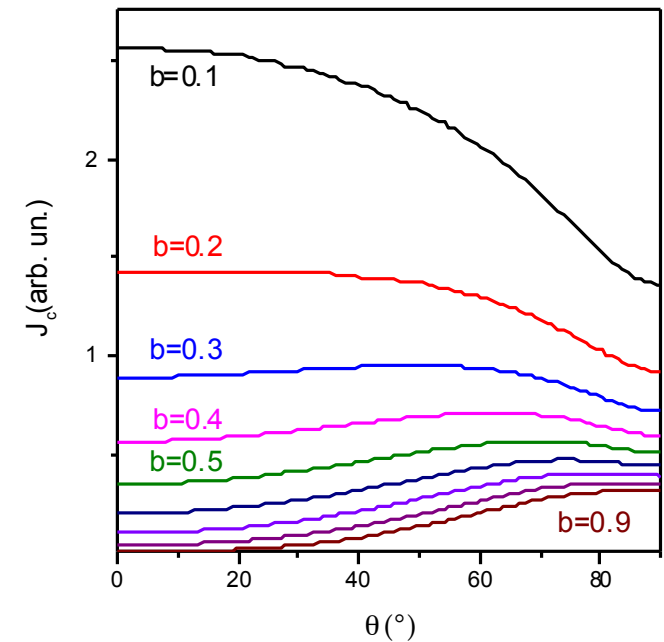


Field scaling

$$J_c(B, \theta) = J_c(B \epsilon(\theta))$$

Peak at 90 °

$\gamma=5$



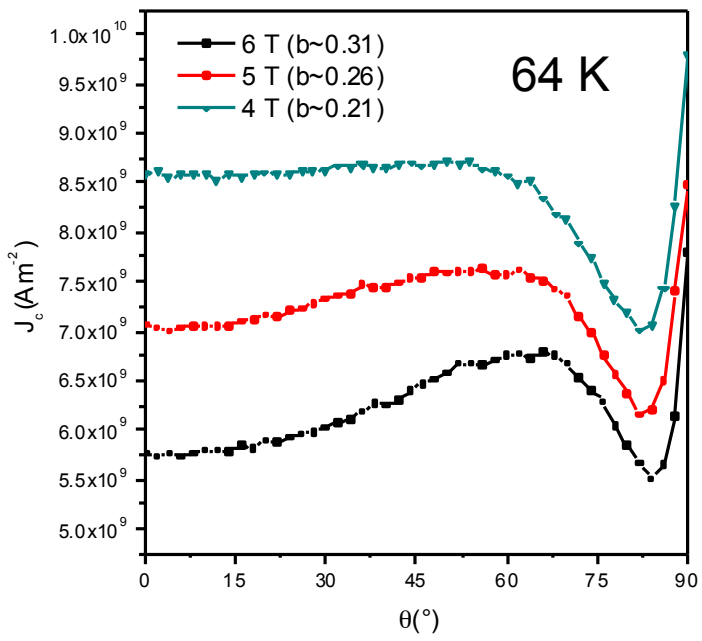
Field and current scaling

$$J_c(B, \theta) = \epsilon(\theta) J_c(B \epsilon(\theta))$$

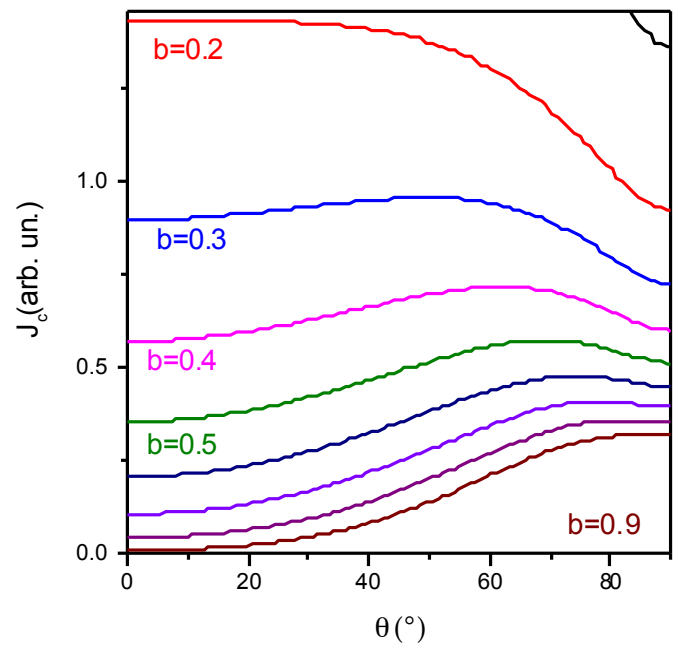
Peak position depends on magnetic field



Angular dependence of J_c



SuperPower SCS 4050



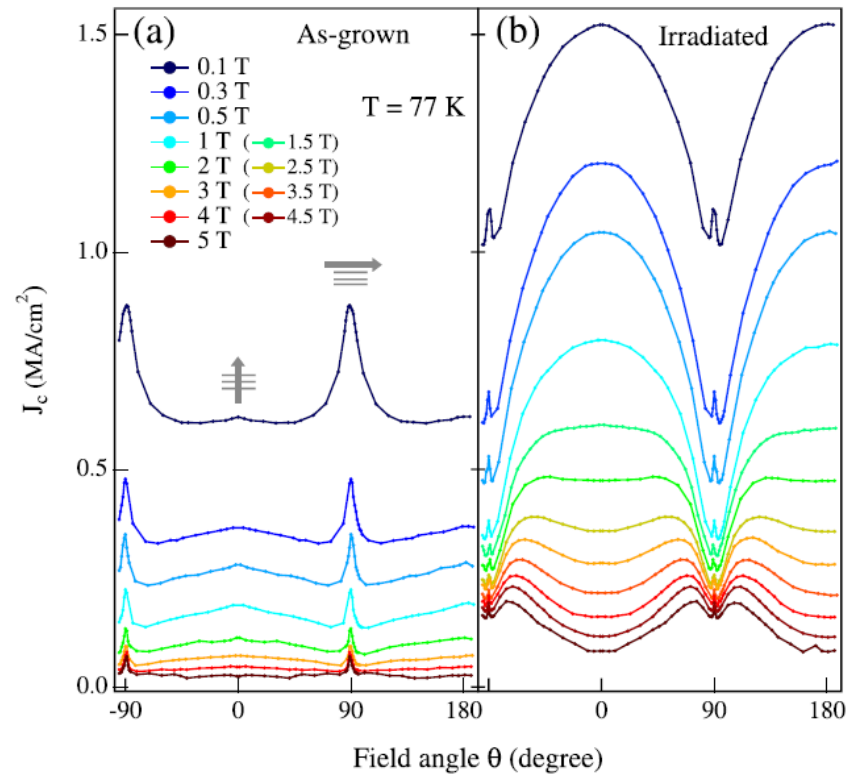
- Qualitative agreement.
- Quantitative agreement unsatisfactorily. (Pre-existing defect structure?)



Data from literature

YBCO film, 3 MeV Au-ions

Model based on flux-lattice shear (Kramer mechanism) leads to the same prediction for $J_c(\theta)$.



H. Matsui et al., J. Appl. Phys 114 (2013) 233911



Conclusions

- Peaks in $J_c(\theta)$ potentially arise from spherical defects with a size comparable to or larger than the coherence length.
- The peak is located at 0° ($H||c$) at low fields and shifts to 90° ($H||ab$) near B_{irr} .
- This qualitative behavior was found in neutron irradiated coated conductors.
- Further work is needed to establish quantitative agreement.

